気象学特論 (ab) (2013 年度秋学期) 最終テスト 解答用紙 (1)

学籍番号: 氏名:

1. (1)

②の両辺をpで偏微分すると、

$$\left(\frac{\partial}{\partial t} + \vec{u}_g \bullet \nabla_p\right) \left\{ \frac{\partial}{\partial p} \left(\frac{f_0}{s^2} \frac{\partial}{\partial p} \Psi_g \right) \right\} + \left(\frac{\partial u_g}{\partial p} \frac{\partial}{\partial x} + \frac{\partial v_g}{\partial p} \frac{\partial}{\partial y} \right) \left(\frac{f_0}{s^2} \frac{\partial}{\partial p} \Psi_g \right) = -\frac{\partial}{\partial p} \omega \qquad \textcircled{2}'$$

左辺の第2項は、

$$\left(\frac{\partial u_{g}}{\partial p} \frac{\partial}{\partial x} + \frac{\partial v_{g}}{\partial p} \frac{\partial}{\partial y}\right) \left(\frac{f_{0}}{s^{2}} \frac{\partial}{\partial p} \Psi_{g}\right) = \left\{-\left(\frac{\partial}{\partial p} \frac{\partial \Psi_{g}}{\partial y}\right) \frac{\partial}{\partial x} + \left(\frac{\partial}{\partial p} \frac{\partial \Psi_{g}}{\partial x}\right) \frac{\partial}{\partial y}\right\} \left(\frac{f_{0}}{s^{2}} \frac{\partial \Psi_{g}}{\partial p}\right) \\
= \frac{f_{0}}{s^{2}} \left\{-\left(\frac{\partial}{\partial y} \frac{\partial \Psi_{g}}{\partial p}\right) \left(\frac{\partial}{\partial x} \frac{\partial \Psi_{g}}{\partial p}\right) + \left(\frac{\partial}{\partial x} \frac{\partial \Psi_{g}}{\partial p}\right) \left(\frac{\partial}{\partial y} \frac{\partial \Psi_{g}}{\partial p}\right)\right\} \\
= 0$$

だから、②'は、

$$\left(\frac{\partial}{\partial t} + \vec{u}_g \bullet \nabla_p\right) \left\{ \frac{\partial}{\partial p} \left(\frac{f_0}{s^2} \frac{\partial}{\partial p} \Psi_g \right) \right\} = -\frac{\partial}{\partial p} \omega$$

$$(2)$$

 $(1) + f_0 \times (2)$ " $\sharp 0$

$$\left(\frac{\partial}{\partial t} + \vec{u}_{g} \bullet \nabla_{p}\right) \left\{ f_{0} + \beta y + \nabla_{p}^{2} \Psi_{g} + \frac{\partial}{\partial p} \left(\frac{f_{0}^{2}}{s^{2}} \frac{\partial}{\partial p} \Psi_{g} \right) \right\} = 0$$

したがって、

$$q = f_0 + \beta y + \nabla_p^2 \Psi_g + \frac{\partial}{\partial p} \left(\frac{f_0^2}{s^2} \frac{\partial}{\partial p} \Psi_g \right)$$

(2)

①の両辺をpで偏微分すると、

$$\frac{\partial}{\partial t} \left\{ \frac{\partial}{\partial p} \left(\nabla_p^2 \Psi_g \right) \right\} + \frac{\partial}{\partial p} \left\{ \vec{\mu}_g \bullet \nabla_p \left(f_0 + \beta y + \nabla_p^2 \Psi_g \right) \right\} = f_0 \frac{\partial^2}{\partial p^2} \omega$$
 ①

②の両辺に ∇_p^2 を作用させると、

$$\frac{\partial}{\partial t} \left\{ \nabla_{p}^{2} \left(\frac{f_{0}}{s^{2}} \frac{\partial}{\partial p} \Psi_{g} \right) \right\} + \nabla_{p}^{2} \left\{ \vec{u}_{g} \bullet \nabla_{p} \left(\frac{f_{0}}{s^{2}} \frac{\partial}{\partial p} \Psi_{g} \right) \right\} = -\nabla_{p}^{2} \omega$$

$$\textcircled{2}^{"}$$

$$\frac{f_0}{s^2} \times (1)' - (2)''' \downarrow \emptyset,$$

$$\left(\nabla_{p}^{2} + \frac{f_{0}^{2}}{s^{2}} \frac{\partial^{2}}{\partial p^{2}}\right) \omega = \frac{f_{0}}{s^{2}} \frac{\partial}{\partial p} \left\{ \vec{u}_{g} \bullet \nabla_{p} \left(f_{0} + \beta y + \nabla_{p}^{2} \Psi_{g}\right) \right\} - \nabla_{p}^{2} \left\{ \vec{u}_{g} \bullet \nabla_{p} \left(\frac{f_{0}}{s^{2}} \frac{\partial}{\partial p} \Psi_{g}\right) \right\} \right\}$$

したがって、

$$\underline{A = f_0 + \beta y + \nabla_p^2 \Psi_g}, \quad \underline{B = \frac{f_0}{s^2} \frac{\partial}{\partial p} \Psi_g}$$

気象学特論 (a b) (2013 年度秋学期) 最終テスト 解答用紙 (2)

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2. (1)

①に②を代入して、

$$(-i\omega + iUk)(-k^2 - l^2)\hat{\Psi}\exp[ik(kx + ly - \omega t)] + i\beta k\hat{\Psi}\exp[ik(kx + ly - \omega t)] = 0$$

両辺を
$$\hat{\Psi}$$
 exp $[ik(kx+ly-\omega t)]$ で割って、
$$(-i\omega+iUk)(-k^2-l^2)+i\beta k=0$$

$$(\omega-Uk)(k^2+l^2)+\beta k=0$$

$$\omega=Uk-\frac{\beta k}{k^2+l^2}$$

(10)

(2)

分散関係式に $\omega = 0$ を代入して、

$$0 = Uk - \frac{\beta k}{k^2 + l^2}$$
$$k^2 + l^2 = \frac{\beta}{U}$$
$$K = \sqrt{\frac{\beta}{U}}$$

$$c_{g_x} = \frac{\partial \omega}{\partial k} = U - \frac{\beta (k^2 + l^2) - 2\beta k^2}{(k^2 + l^2)^2} = U + \frac{\beta (k^2 - l^2)}{(k^2 + l^2)^2}$$

(2) の結果より、 $\beta = U(k^2 + l^2)$ だから、

$$c_{gx} = U + \frac{k^2 - l^2}{k^2 + l^2}U = \frac{2k^2}{\underline{k^2 + l^2}}U$$

(10)

(4)

$$c_{gy} = \frac{\partial \omega}{\partial l} = \frac{2\beta kl}{\left(k^2 + l^2\right)^2}$$

(2) の結果より、 $\beta = U(k^2 + l^2)$ だから、

$$c_{gy} = \frac{2kl(k^2 + l^2)}{(k^2 + l^2)^2}U = \frac{2kl}{\underline{k^2 + l^2}}U$$

(10)

(5)

$$c_{gx} - U = \left(\frac{2k^2}{k^2 + l^2} - 1\right)U = \frac{k^2 - l^2}{k^2 + l^2}U$$
$$c_{gy} = \frac{2kl}{k^2 + l^2}U$$

だから、

$$(c_{gx} - U)^{2} + c_{gy}^{2} = \frac{(k^{2} - l^{2})^{2}}{(k^{2} + l^{2})^{2}} U^{2} + \frac{4k^{2}l^{2}}{(k^{2} + l^{2})^{2}} U^{2} = U^{2}$$
$$\underline{(c_{gx} - U)^{2} + c_{gy}^{2}} = U^{2}$$

(10)

気象学特論 (a b) (2013 年度秋学期) 最終テスト 解答用紙 (3)

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3. (1)

④を②に代入すると、

$$\frac{-k^2(-i\nu - ikU)\hat{\Psi}_3 = \hat{\omega}_2}{2}$$

④を③に代入すると、

$$-i\nu(\hat{\Psi}_{1} - \hat{\Psi}_{3}) - ikU(\hat{\Psi}_{1} + \hat{\Psi}_{3}) = \frac{1}{\lambda^{2}}\hat{\omega}_{2}$$

$$(-i\nu - ikU)\hat{\Psi}_{1} + (i\nu - ikU)\hat{\Psi}_{3} = \frac{1}{\lambda^{2}}\hat{\omega}_{2}$$

$$(3)$$

(10)

(2)

連立方程式①'~③'は、

$$\begin{pmatrix}
-k^2(-i\nu+ikU) & 0 & -1 \\
0 & -k^2(-i\nu-ikU) & 1 \\
-i\nu-ikU & i\nu-ikU & -\frac{1}{\lambda^2}
\end{pmatrix} \begin{pmatrix} \hat{\Psi}_1 \\ \hat{\Psi}_3 \\ \hat{\omega}_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$$

と書くことができる。したがって、

$$A = \begin{pmatrix} -k^{2}(-iv + ikU) & 0 & -1\\ 0 & -k^{2}(-iv - ikU) & 1\\ -iv - ikU & iv - ikU & -\frac{1}{\lambda^{2}} \end{pmatrix}$$

(3)

行列 A の行列式がゼロだから、

$$k^{4}(-i\nu+ikU)(-i\nu-ikU)\left(-\frac{1}{\lambda^{2}}\right) - (-k^{2})(-i\nu+ikU)(i\nu-ikU)$$

$$-(-1)(-i\nu-ikU)(-k^{2})(-i\nu-ikU) = 0$$

$$-k^{4}(k^{2}U^{2}-\nu^{2})\frac{1}{\lambda^{2}} + k^{2}(\nu-kU)(\nu-kU) + k^{2}(\nu+kU)(\nu+kU) = 0$$

$$-k^{2}(k^{2}U^{2}-\nu^{2})\frac{1}{\lambda^{2}} + 2(\nu^{2}+k^{2}U^{2}) = 0$$

$$\left(\frac{1}{\lambda^{2}}k^{2} + 2\right)\nu^{2} - k^{2}U^{2}\left(\frac{1}{\lambda^{2}}k^{2} - 2\right) = 0$$

$$(k^{2}+2\lambda^{2})\nu^{2} - k^{2}U^{2}(k^{2}-2\lambda^{2}) = 0$$

$$\nu^{2} = \frac{k^{2}U^{2}(k^{2}-2\lambda^{2})}{k^{2}+2\lambda^{2}}$$